

# Measurement of

$$\Gamma_{ee}(J/\psi) \cdot \mathcal{B}(J/\psi \rightarrow e^+e^-) \text{ and } \Gamma_{ee}(J/\psi) \cdot \mathcal{B}(J/\psi \rightarrow \mu^+\mu^-)$$

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## Abstract

The products of the electron width of the  $J/\psi$  meson and the branching fraction of its decays to the lepton pairs were measured using data from the KEDR experiment at the VEPP-4M electron-positron collider. The results are

$$\Gamma_{ee} \times \Gamma_{ee}/\Gamma = 0.3323 \pm 0.0064 (\text{stat.}) \pm 0.0048 (\text{syst.}) \text{ keV},$$

$$\Gamma_{ee} \times \Gamma_{\mu\mu}/\Gamma = 0.3318 \pm 0.0052 (\text{stat.}) \pm 0.0063 (\text{syst.}) \text{ keV},$$

$$\Gamma_{ee} \times (\Gamma_{ee} + \Gamma_{\mu\mu})/\Gamma = 0.6641 \pm 0.0082 (\text{stat.}) \pm 0.0100 (\text{syst.}) \text{ keV},$$

$$\Gamma_{ee}/\Gamma_{\mu\mu} = 1.002 \pm 0.021 (\text{stat.}) \pm 0.013 (\text{syst.})$$

and can be used to improve the accuracy of the leptonic and full widths and test leptonic universality. Assuming  $e\mu$  universality and using the world average value of the leptonic branching fraction, we also determine the leptonic  $\Gamma_{\ell\ell} = 5.59 \pm 0.12 \text{ keV}$  and total  $\Gamma = 94.1 \pm 2.7 \text{ keV}$  widths of the  $J/\psi$  meson. V. V. Anashin, et al., Phys. Lett. B 685, 134 (2010), arXiv:0912.1082.

## 1. VEPP-4M collider and KEDR detector

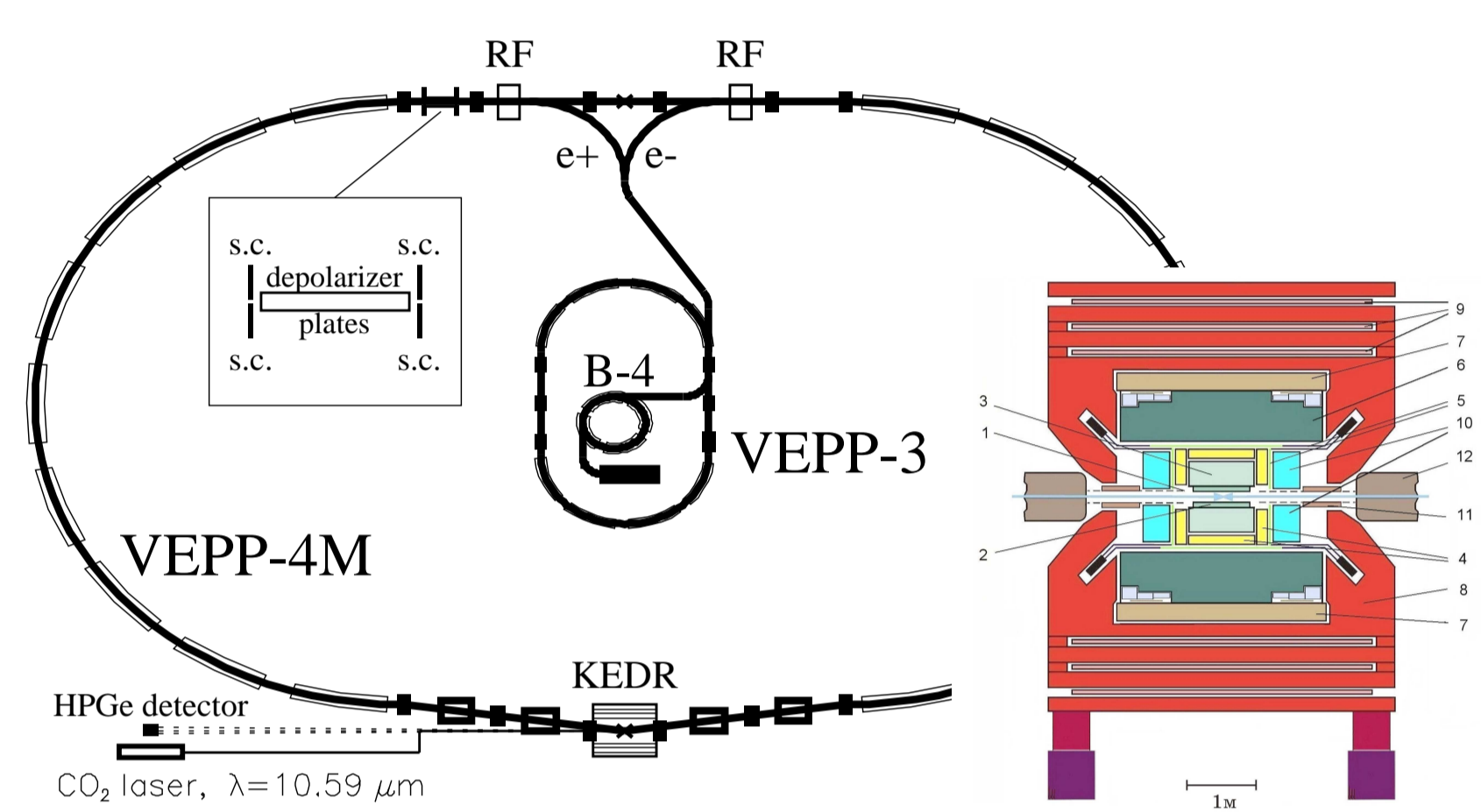


Figure 1: VEPP-4M/KEDR complex with the resonant depolarization and the infrared light Compton backscattering facilities.

One of the main features of the VEPP-4M is a possibility of precise energy determination. Between calibrations the energy interpolation in the  $J/\psi$  energy range has the accuracy of  $6 \cdot 10^{-6}$  ( $\approx 10 \text{ keV}$ ).

## 2. Experiment description

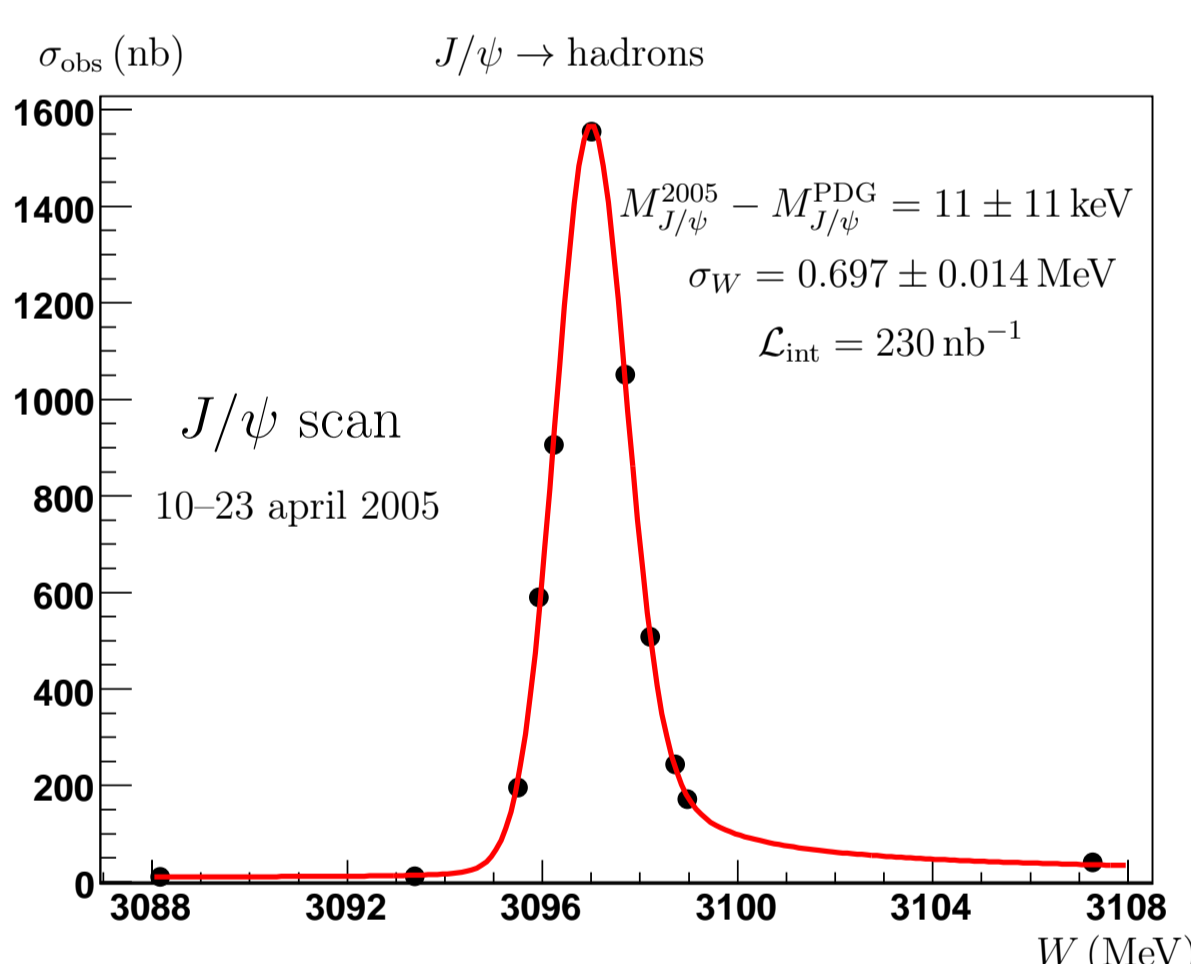


Figure 2: Observed cross section of  $e^+e^- \rightarrow \text{hadrons}$ .

A data sample used for this analysis comprises 230  $\text{nb}^{-1}$  collected at 11 energy points in the  $J/\psi$  energy range. This corresponds to approximately 15000  $J/\psi \rightarrow e^+e^-$  decays. During this scan, 26 calibrations of the beam energy have been done using resonant depolarization.

## 3. Theoretical $e^+e^- \rightarrow \ell^+\ell^-$ cross section

The analytical expressions for the cross section of the process  $e^+e^- \rightarrow \ell^+\ell^-$  with radiative corrections taken into account in the soft photon approximation were first derived by Ya. A. Azimov et al., JETP Lett. 21, 172 (1975). With some up-to-day modifications one obtains in the vicinity of a narrow resonance:

$$\left(\frac{d\sigma}{d\Omega}\right)^{ee \rightarrow ee} \approx \frac{1}{M^2} \left\{ \frac{9\Gamma_{ee}^2}{4\Gamma M} (1 + \cos^2\theta) (1 + \delta_{\text{sf}}) \text{Im } \mathcal{F} - \frac{3\alpha\Gamma_{ee}}{2M} \left[ (1 + \cos^2\theta) - \frac{(1 + \cos\theta)^2}{(1 - \cos\theta)} \right] \text{Re } \mathcal{F} \right\} + \left(\frac{d\sigma}{d\Omega}\right)_{\text{QED}}^{ee}$$

$$\left(\frac{d\sigma}{d\Omega}\right)^{ee \rightarrow \mu\mu} \approx \frac{3}{4M^2} (1 + \delta_{\text{sf}}) (1 + \cos^2\theta) \times \left\{ \frac{3\Gamma_{ee}\Gamma_{\mu\mu}}{\Gamma M} \text{Im } \mathcal{F} - \frac{2\alpha\sqrt{\Gamma_{ee}\Gamma_{\mu\mu}}}{M} \text{Re } \frac{\mathcal{F}}{1 - \Pi_0} \right\} + \left(\frac{d\sigma}{d\Omega}\right)_{\text{QED}}^{\mu\mu}$$

where a correction  $\delta_{\text{sf}}$  follows from the structure function approach of E. A. Kuraev and V. S. Fadin, Sov. J. Nucl. Phys. 41, 466 (1985).

$$\mathcal{F} = \frac{\pi\beta}{\sin\pi\beta} \left( \frac{M/2}{-W + M - i\Gamma/2} \right)^{1-\beta}, \quad \beta = \frac{4\alpha}{\pi} \left( \ln \frac{W}{m_e} - \frac{1}{2} \right)$$

## 4. Data analysis

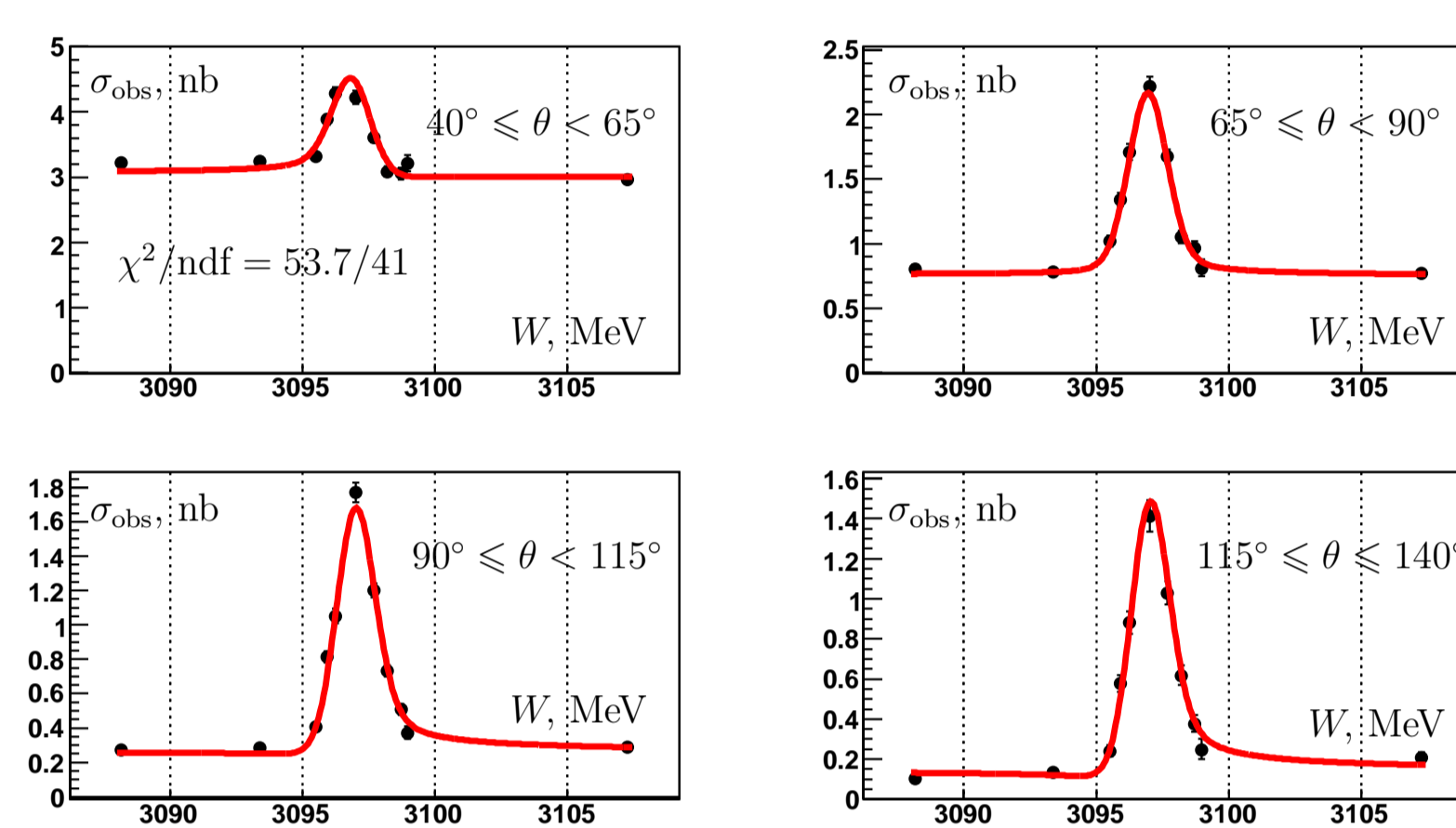


Figure 3: Fits to experimental data for  $e^+e^- \rightarrow e^+e^-$ .

At the  $i$ -th energy point  $E_i$  and the  $j$ -th angular interval  $\theta_j$ , the expected number of  $e^+e^- \rightarrow e^+e^-$  events was parameterized as

$$N_{\text{exp}}(E_i, \theta_j) = \mathcal{R}_{\mathcal{L}} \times \mathcal{L}(E_i) \times \left( \sigma_{\text{res}}^{\text{theor}}(E_i, \theta_j) \cdot \varepsilon_{\text{res}}^{\text{sim}}(E_i, \theta_j) + \sigma_{\text{inter}}^{\text{theor}}(E_i, \theta_j) \cdot \varepsilon_{\text{inter}}^{\text{sim}}(E_i, \theta_j) + \sigma_{\text{Bhabha}}^{\text{sim}}(E_i, \theta_j) \cdot \varepsilon_{\text{Bhabha}}^{\text{sim}}(E_i, \theta_j) \right)$$

where  $\mathcal{L}(E_i)$  — the integrated luminosity measured by the luminosity monitor at the  $i$ -th energy point;  $\sigma^{\text{theor}}$  — the theoretical cross sections for resonance, interference and Bhabha contributions,  $\varepsilon^{\text{sim}}$  — the detector efficiencies obtained from simulation.

In this formula the following free parameters were used: the product  $\Gamma_{ee} \times \Gamma_{ee}/\Gamma$ , which determines the magnitude of the resonance signal; the electron width  $\Gamma_{ee}$ , which specifies the amplitude of the interference wave; the coefficient  $\mathcal{R}_{\mathcal{L}}$ , which provides the absolute calibration of the luminosity monitor.

The dominant uncertainty of the  $\Gamma_{ee} \times \Gamma_{ee}/\Gamma$  result is associated with the luminosity monitor instability.

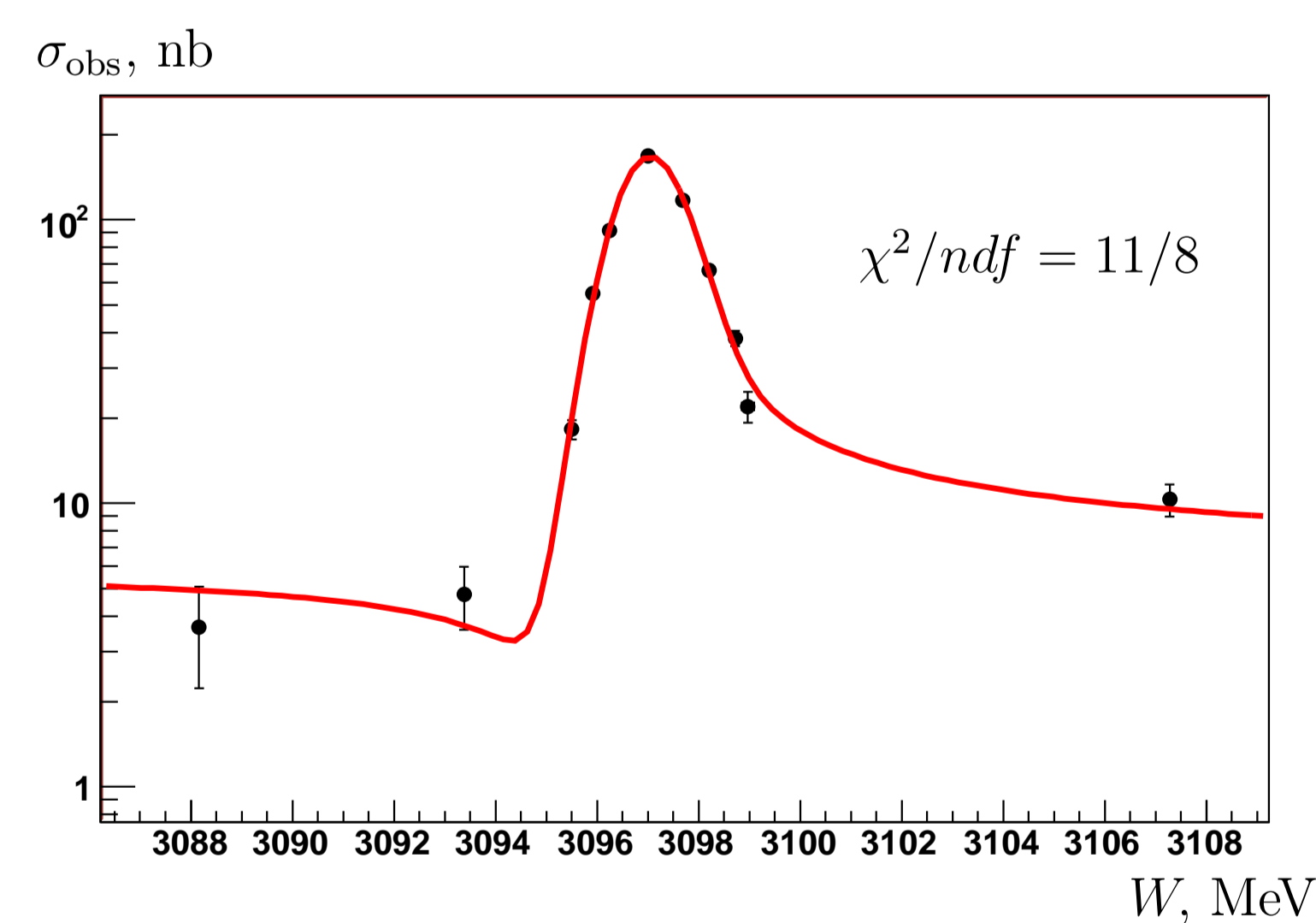


Figure 4: Fit to experimental data for  $e^+e^- \rightarrow \mu^+\mu^-$ .

The expected number of  $e^+e^- \rightarrow \mu^+\mu^-$  events was parameterized in the form:

$$N_{\text{exp}}(E_i) = \mathcal{R}_{\mathcal{L}} \times \mathcal{L}(E_i) \times \left( \sigma_{\text{res}}^{\text{theor}}(E_i) \cdot \varepsilon_{\text{res}}^{\text{sim}}(E_i) + \sigma_{\text{inter}}^{\text{theor}}(E_i) \cdot \varepsilon_{\text{inter}}^{\text{sim}}(E_i) + \sigma_{\text{bg}}^{\text{theor}}(E_i) \cdot \varepsilon_{\text{bg}}^{\text{sim}}(E_i) + F_{\text{cosmic}} \times T_i \right)$$

with the same meaning of  $\mathcal{R}_{\mathcal{L}}$  and  $\mathcal{L}(E_i)$  as for  $e^+e^- \rightarrow e^+e^-$ .  $\mathcal{R}_{\mathcal{L}}$  was fixed from the  $e^+e^- \rightarrow e^+e^-$  fit and  $T_i$  is the live data taking time.

The following free parameters were used: the product  $\Gamma_{ee} \times \Gamma_{\mu\mu}/\Gamma$ , which determines the magnitude of the resonance signal; the square root of electron and muon widths  $\sqrt{\Gamma_{ee}\Gamma_{\mu\mu}}$ , which specifies the amplitude of the interference wave; the rate of cosmic events,  $F_{\text{cosmic}}$ , that passed the selection criteria for the  $e^+e^- \rightarrow \mu^+\mu^-$  events.

The dominant uncertainty of the  $\Gamma_{ee} \times \Gamma_{\mu\mu}/\Gamma$  result is associated with the absolute luminosity calibration done in the  $e^+e^-$ -channel.

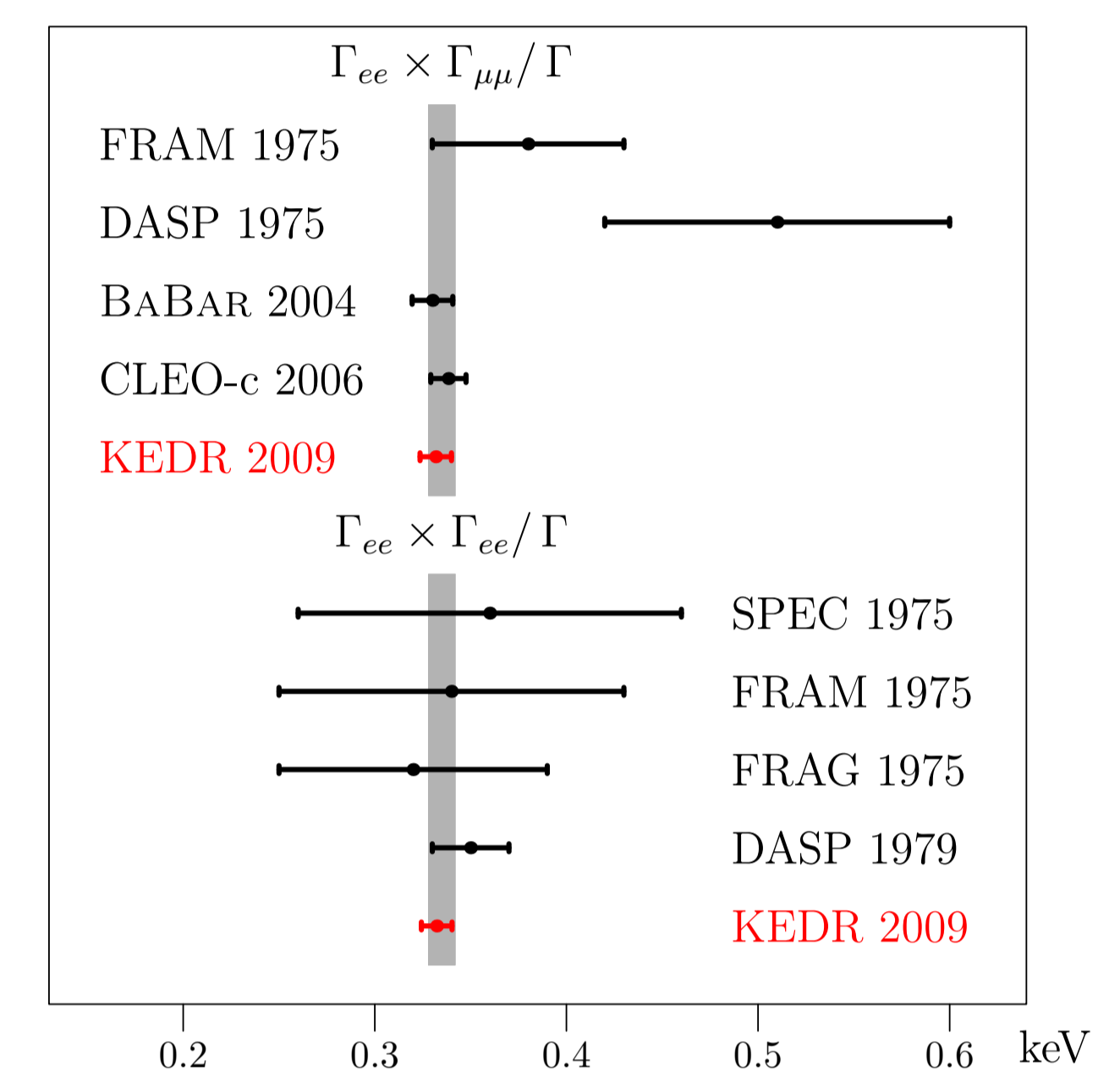


Figure 5:  $\Gamma_{ee} \times \Gamma_{ee}/\Gamma$  and  $\Gamma_{ee} \times \Gamma_{\mu\mu}/\Gamma$  comparison

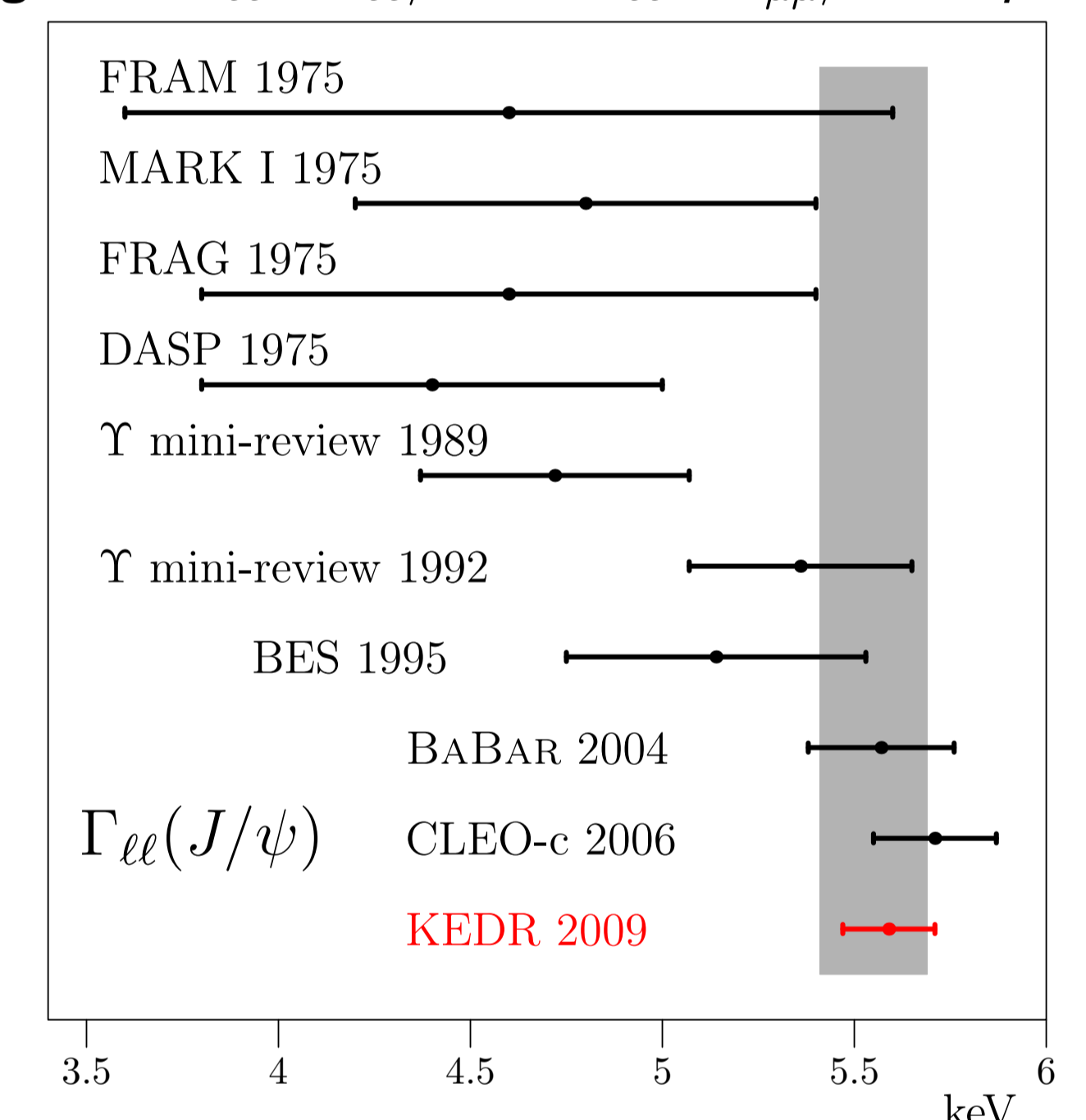


Figure 6:  $\Gamma_{\ell\ell}$  comparison

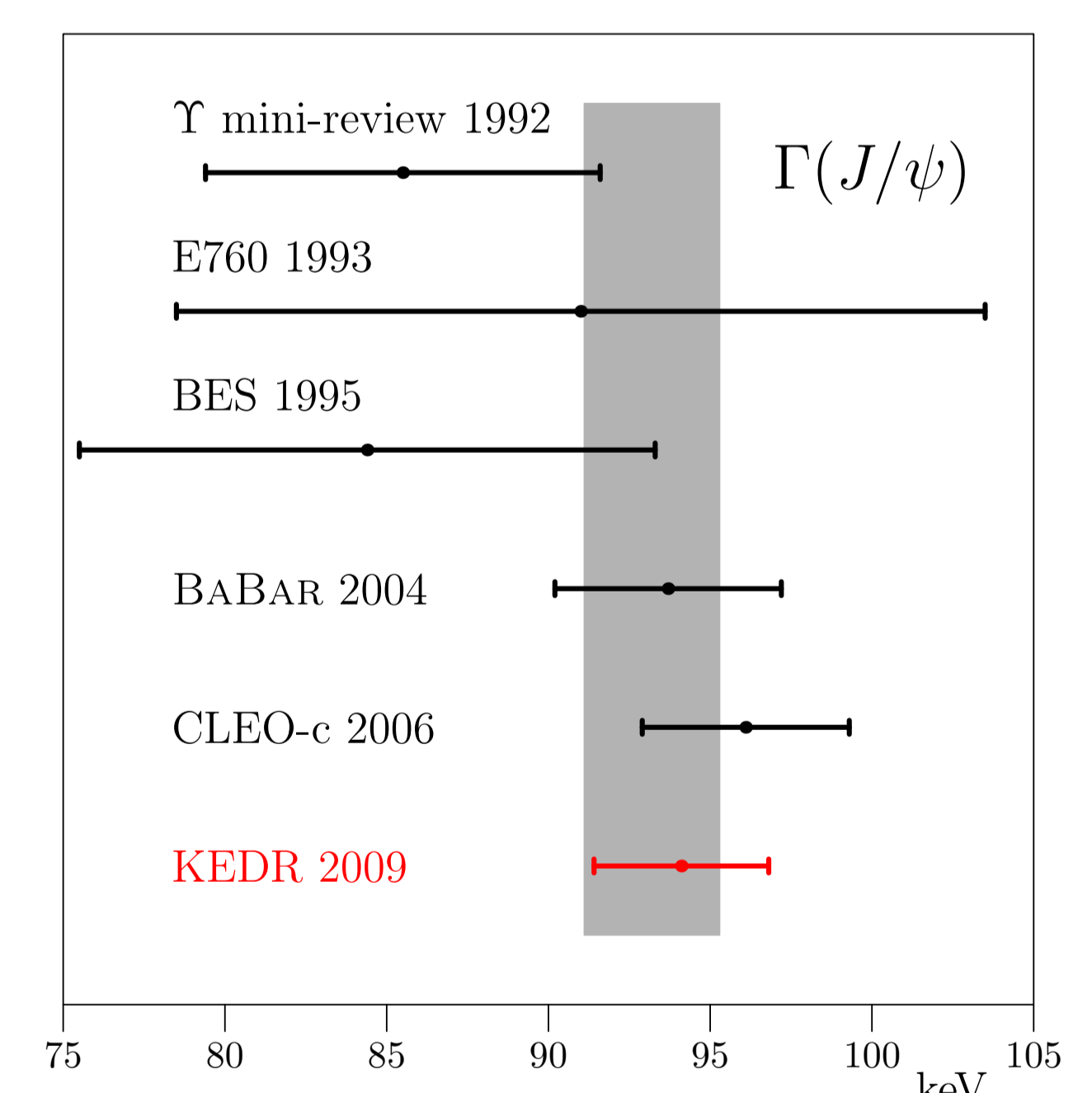


Figure 7:  $\Gamma$  comparison

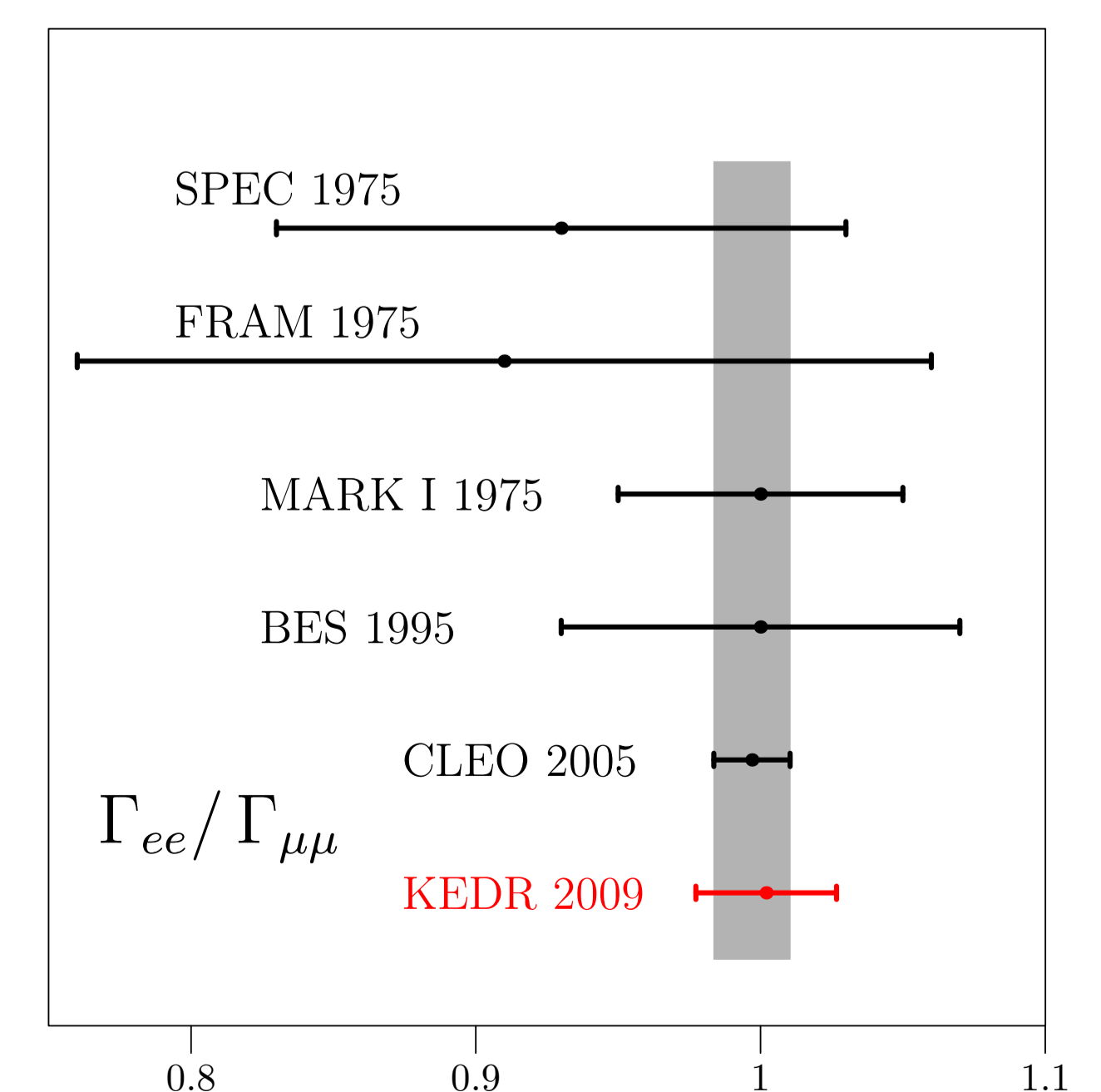


Figure 8:  $\Gamma_{ee}/\Gamma_{\mu\mu}$  comparison